

## 593.

## A SHEEPHANKS' PROBLEM (1866).

[From the *Messenger of Mathematics*, vol. iv. (1875), pp. 34—36.]

APPLY the formulæ of elliptic motion to determine the motion of a body let fall from the top of a tower at the equator.

The earth is regarded as rotating with the angular velocity  $\omega$  round a fixed axis, so that the body is in fact projected from the apocentre with an angular velocity  $= \omega$ ; and we write  $\alpha$  for the equatorial radius,  $\beta$  for the height of the tower; then  $g$  denoting the force of gravity, and  $\mu, h, n, a, e, \theta$ , as in the theory of elliptic motion, we have

$$\mu = n^2 a^3 = g \alpha^2,$$

$$h = (\alpha + \beta)^2 \omega = n a^2 \sqrt{1 - e^2},$$

$$\alpha + \beta = a(1 + e);$$

whence

$$(\alpha + \beta)^4 \omega^2 = g \alpha^2 a (1 - e^2),$$

$$(\alpha + \beta) = a(1 + e),$$

$$\frac{(\alpha + \beta)^3 \omega^2}{g \alpha^2} = \frac{\alpha \omega^2}{g} \left(1 + \frac{\beta}{\alpha}\right)^3 = 1 - e,$$

where  $\frac{\alpha \omega^2}{g}$  = ratio of centrifugal force to gravity,

$$= \frac{1}{288},$$

so that  $1 - e$  is small;

$$r = \frac{a(1 - e^2)}{1 - e \cos \theta}, = \frac{(\alpha + \beta)(1 - e)}{1 - e \cos \theta},$$

whence

$$1 - e \cos \theta = \frac{(\alpha + \beta)(1 - e)}{r}.$$

Suppose

$$r = \alpha,$$

$$1 - e \cos \theta = (1 - e) \left( 1 + \frac{\beta}{\alpha} \right) = 1 - e + \frac{\beta}{\alpha} (1 - e),$$

which is nearly

$$= 1 - e,$$

that is,  $\theta$  is small, and therefore approximately

$$1 - e + \frac{1}{2} e \theta^2 = 1 - e + \frac{\beta}{\alpha} (1 - e),$$

or

$$\theta^2 = \frac{2\beta}{\alpha} \frac{1 - e}{e};$$

we then have

$$\begin{aligned} r^2 d\theta = h dt, \text{ or } dt &= \frac{r^2 d\theta}{h} = \frac{(\alpha + \beta)^2 (1 - e)^2}{(\alpha + \beta)^2 \omega (1 - e \cos \theta)^2} d\theta \\ &= \frac{(1 - e)^2}{\omega^2 (1 - e \cos \theta)^2} d\theta \\ &= \frac{(1 - e)^2}{\omega (1 - e + \frac{1}{2} e \theta^2)^2} d\theta \\ &= \frac{1}{\omega \left( 1 + \frac{\frac{1}{2} e}{1 - e} \theta^2 \right)^2} d\theta \\ &= \frac{1}{\omega} \left( 1 - \frac{e}{1 - e} \theta^2 \right) d\theta; \end{aligned}$$

that is,

$$\omega dt = \left( 1 - \frac{e}{1 - e} \theta^2 \right) d\theta.$$

Integrating, we have

$$\theta \left( 1 - \frac{\frac{1}{3} e \theta^2}{1 - e} \right) = \omega t,$$

where  $\omega t$  = earth's rotation in time  $t$ , =  $\phi$  suppose; therefore

$$\theta \left( 1 - \frac{\frac{1}{3} e \theta^2}{1 - e} \right) = \phi,$$

hence, if  $\theta$  be as above, the angle described in falling to the surface,

$$\frac{e \theta^2}{1 - e} = \frac{2\beta}{\alpha},$$

$$\theta \left( 1 - \frac{2}{3} \frac{\beta}{\alpha} \right) = \phi,$$

$$\theta - \phi = \frac{2}{3} \frac{\beta}{\alpha} \theta = \frac{2}{3} \frac{\beta}{\alpha} \sqrt{\left( \frac{2\beta}{\alpha} \frac{1 - e}{e} \right)}.$$

Writing herein

$$\frac{1-e}{e}, = 1-e, = \frac{\alpha\omega^2}{g},$$

this is

$$= \frac{\beta}{\alpha} \sqrt{\left(\frac{2\beta}{\alpha} \frac{\alpha\omega^2}{g}\right)} = \frac{(2\beta)^{\frac{3}{2}}}{3\alpha^{\frac{3}{2}}} \sqrt{\left(\frac{\alpha\omega^2}{g}\right)},$$

viz.

$$\theta - \phi = \frac{2^{\frac{3}{2}} \cdot \beta^{\frac{3}{2}}}{3\alpha^{\frac{3}{2}}} \sqrt{\left(\frac{\alpha\omega^2}{g}\right)},$$

whence

$$\alpha(\theta - \phi) = \frac{2\sqrt{(2)}}{3} \sqrt{\left(\frac{\beta}{\alpha}\right)} \sqrt{\left(\frac{\alpha\omega^2}{g}\right)} \beta,$$

or say

$$\frac{\alpha(\theta - \phi)}{\beta} = \frac{2\sqrt{(2)}}{3} \sqrt{\left(\frac{\beta}{\alpha}\right)} \sqrt{\left(\frac{\alpha\omega^2}{g}\right)},$$

where  $\alpha(\theta - \phi)$  is the distance at which the body falls from the foot of the tower.

Substituting for  $\sqrt{\left(\frac{\alpha\omega^2}{g}\right)}$  its value,  $= \frac{1}{17}$ , we have

$$\frac{\alpha(\theta - \phi)}{\beta} = \frac{2\sqrt{(2)}}{51} \sqrt{\left(\frac{\beta}{\alpha}\right)}, = 0.056 \sqrt{\left(\frac{\beta}{\alpha}\right)}.$$