

466.

SECOND NOTE ON THE LUNAR THEORY.

[From the *Monthly Notices of the Royal Astronomical Society*, vol. xxv. (1864—1865), pp. 203—207.]

THE elliptic values of

r , the radius vector,

v , the longitude,

y , the latitude,

are functions of

a , the mean distance,

e , the excentricity,

γ , the tangent of the inclination,

l , the mean longitude,

c , the mean anomaly,

g , the mean distance from node;

see my Note in the last *Monthly Notice*, p. 182, [465], where, for the present purpose, $\frac{a}{r}$ should be written instead of $\frac{1}{r}$; and it is there shown that the disturbed values, attending only to the coefficients independent of m , are obtained by affecting a, e, γ, c, g, l , with the inequalities

$$\begin{aligned} \delta a &= 0 \\ \delta e &= -\frac{5}{8} \gamma^2 e \quad \cos \quad 2c - 2g \\ \delta \gamma &= +\frac{5}{8} \gamma e^2 \quad \text{,,} \quad 2c - 2g \\ \delta c &= +\frac{5}{8} \gamma^2 \quad \sin \quad 2c - 2g \\ \delta g &= +\frac{5}{8} e^2 \quad \text{,,} \quad 2c - 2g \\ \delta l &= -\frac{5}{16} \gamma^2 e^2 \quad \text{,,} \quad 2c - 2g, \end{aligned}$$

or, what is the same thing, adding to the elliptic values the inequalities

$$\begin{aligned}
 \delta \frac{a}{r} &= -\frac{5}{4} \gamma^2 e^2 & \cos & 2g \\
 & -\frac{5}{8} \gamma^2 e & & c - 2g, \\
 \delta v &= -\frac{5}{16} \gamma^2 e^2 & \sin & 2c \\
 & -\frac{25}{16} \gamma^2 e^2 & & 2g \\
 & +\frac{5}{4} \gamma^2 e & & c - 2g \\
 & -\frac{5}{16} \gamma^2 e^2 & & 2c - 2g, \\
 \delta y &= -\frac{5}{8} \gamma e^3 + \frac{5}{8} \gamma^3 e & \sin & c - g \\
 & +\frac{5}{8} \gamma e^2 & & 2c - g \\
 & +\frac{5}{8} \gamma e^3 & & 3c - g \\
 & +\frac{5}{8} \gamma^3 e & & c - 3g.
 \end{aligned}$$

I propose to show how these results may be obtained by the method of the variation of the elements. For this purpose, treating a, e, γ, c, g, l , as elements, the proper formulæ are obtained very readily from those given in my "Memoir on the Problem of Disturbed Elliptic Motion," *Mem. R. Ast. Soc.*, vol. xxvii. (1859), pp. 1—29, [212]; viz., writing c in place of g , the formulæ, p. 25, give the variations of $a, e, c, \mathcal{T}, \theta, \phi$; we have then

$$\begin{aligned}
 g &= c + \mathcal{T} \\
 l &= c + \mathcal{T} + \theta \\
 \gamma &= \tan \phi,
 \end{aligned}$$

and therefore

$$\begin{aligned}
 dg &= dc + d\mathcal{T} \\
 dl &= dc + d\mathcal{T} + d\theta \\
 d\gamma &= (1 + \gamma^2) d\phi,
 \end{aligned}$$

which give for the transformation of the differential coefficients of Ω ,

$$\begin{aligned}
 \frac{d\Omega}{dc} &= \frac{d\Omega}{dc} + \frac{d\Omega}{dg} + \frac{d\Omega}{dl} \\
 \frac{d\Omega}{d\mathcal{T}} &= \frac{d\Omega}{dg} + \frac{d\Omega}{dl} \\
 \frac{d\Omega}{d\theta} &= \frac{d\Omega}{dl} \\
 \frac{d\Omega}{d\phi} &= (1 + \gamma^2) \frac{d\Omega}{d\gamma}
 \end{aligned}$$

and the formulæ finally become

$$\begin{aligned}
 \frac{da}{dt} &= \frac{2}{na} \frac{d\Omega}{dc} + \frac{2}{na} \frac{d\Omega}{dg} + \frac{2}{na} \frac{d\Omega}{dl}, \\
 \frac{de}{dt} &= \frac{1 - e^2}{na^2 e} \frac{d\Omega}{dc} + \frac{1 - e^2 - \sqrt{1 - e^2}}{na^2 e} \frac{d\Omega}{dg} + \frac{1 - e^2 - \sqrt{1 - e^2}}{na^2 e} \frac{d\Omega}{dl},
 \end{aligned}$$

$$\frac{d\gamma}{dt} = \frac{1 + \gamma^2}{na^2 \sqrt{1 - e^2} \gamma} \frac{d\Omega}{dg} + \frac{(1 + \gamma^2)(1 - \sqrt{1 + \gamma^2})}{na^2 \sqrt{1 - e^2} \gamma} \frac{d\Omega}{dl},$$

$$\frac{dc}{dt} = -\frac{2}{na} \frac{d\Omega}{da} - \frac{1 - e^2}{na^2 e} \frac{d\Omega}{de},$$

$$\frac{dg}{dt} = -\frac{2}{na} \frac{d\Omega}{da} - \frac{1 - e^2 - \sqrt{1 - e^2}}{na^2 e} \frac{d\Omega}{de} - \frac{1 + \gamma^2}{na^2 \sqrt{1 - e^2} \gamma} \frac{d\Omega}{d\gamma},$$

$$\frac{dl}{dt} = -\frac{2}{na} \frac{d\Omega}{da} - \frac{1 - e^2 - \sqrt{1 + e^2}}{na^2 e} \frac{d\Omega}{de} - \frac{(1 + \gamma^2)(1 - \sqrt{1 + \gamma^2})}{na^2 \sqrt{1 - e^2} \gamma} \frac{d\Omega}{d\gamma}.$$

The disturbing function contains the term

$$m^2 n^2 a^2 \left(+ \frac{15}{8} e^2 \gamma^2 \right) \cos 2c - 2g.$$

If after the differentiations we write for greater simplicity $a = 1, n = 1$, we have

$$\frac{d\Omega}{da} = + \frac{15}{8} m^2 e^2 \gamma^2 \cos 2c - 2g,$$

$$\frac{d\Omega}{de} = + \frac{15}{8} m^2 e \gamma^2 \quad ,, \quad 2c - 2g,$$

$$\frac{d\Omega}{d\gamma} = + \frac{15}{8} m^2 e^2 \gamma \quad ,, \quad 2c - 2g,$$

$$\frac{d\Omega}{dc} = - \frac{15}{8} m^2 e^2 \gamma^2 \sin 2c - 2g,$$

$$\frac{d\Omega}{dg} = - \frac{15}{8} m^2 e^2 \gamma^2 \quad ,, \quad 2c - 2g,$$

$$\frac{d\Omega}{dt} = 0,$$

and the formulæ for the variations give

$$\frac{da}{dt} = 2 \left(\frac{d\Omega}{dc} + \frac{d\Omega}{dg} \right) = 0$$

$$\frac{de}{dt} = \frac{1}{e} \frac{d\Omega}{dc} = - \frac{15}{8} m^2 e \gamma^2 \sin 2c - 2g,$$

$$\frac{d\gamma}{dt} = \frac{1}{\gamma} \frac{d\Omega}{dg} = - \frac{15}{8} m^2 e^2 \gamma \quad ,, \quad 2c - 2g,$$

$$\frac{dc}{dt} = -\frac{1}{e} \frac{d\Omega}{de} = - \frac{15}{8} m^2 \gamma^2 \cos 2c - 2g,$$

$$\frac{dg}{dt} = -\frac{1}{\gamma} \frac{d\Omega}{d\gamma} = - \frac{15}{8} m^2 e^2 \quad ,, \quad 2c - 2g,$$

$$\frac{dl}{dt} = -2 \frac{d\Omega}{da} + \frac{1}{2} e \frac{d\Omega}{de} + \frac{1}{2} \gamma \frac{d\Omega}{d\gamma} = \left(-\frac{15}{4} + \frac{15}{8} + \frac{15}{8} \right) - \frac{15}{8} m^2 e^2 \gamma^2 \quad ,, \quad 2c - 2g,$$

but this value of $\frac{dl}{dt}$ is, as will presently be seen, incomplete.

Writing $a + \delta a$, $e + \delta e$, &c., in place of a , e , &c., and observing that the divisor for the integration of the term in $2c - 2g$ is $2(c - g)$, $= -3m^2$, the first five equations give respectively

$$\begin{aligned}\delta a &= 0, \\ \delta e &= -\frac{5}{8} \gamma^2 e \quad \cos \quad 2c - 2g, \\ \delta \gamma &= +\frac{5}{8} \gamma e^2 \quad \text{,,} \quad 2c - 2g, \\ \delta c &= +\frac{5}{8} \gamma^2 \quad \sin \quad 2c - 2g, \\ \delta g &= +\frac{5}{8} e^2 \quad \text{,,} \quad 2c - 2g.\end{aligned}$$

The constant term in Ω is

$$= m^2 n^2 a^2 \left(\frac{1}{4} + \frac{3}{8} e^2 - \frac{3}{8} \gamma^2 \right),$$

and this gives in

$$\frac{dl}{dt} = -2 \frac{d\Omega}{da} + \frac{1}{2} e \frac{d\Omega}{de} + \frac{1}{2} \gamma \frac{d\Omega}{d\gamma},$$

a term

$$\begin{aligned}m^2 \left(-1 - \frac{3}{2} e^2 + \frac{3}{2} \gamma^2 \right. \\ \left. + \frac{3}{8} e^2 - \frac{3}{8} \gamma^2 \right),\end{aligned}$$

which is

$$= m^2 \left(-1 - \frac{9}{8} e^2 + \frac{9}{8} \gamma^2 \right).$$

Substituting for e , γ , their correct values $e + \delta e$, $\gamma + \delta \gamma$, it appears that $\frac{dl}{dt}$ contains the term

$$m^2 \left(-\frac{9}{4} e \delta e + \frac{9}{4} \gamma \delta \gamma \right),$$

which is

$$\begin{aligned}&= m^2 \left(\frac{45}{32} + \frac{45}{32} \right) \frac{45}{16} e^2 \gamma^2 \cos \quad 2c - 2g, \\ &= \frac{45}{16} m^2 e^2 \gamma^2 \quad \text{,,} \quad 2c - 2g,\end{aligned}$$

and joining to this the before-mentioned term

$$= -\frac{15}{8} m^2 e^2 \gamma^2 \quad \text{,,} \quad 2c - 2g,$$

we find

$$\frac{dl}{dt} = \left(\frac{45}{16} - \frac{15}{8} \right) \frac{15}{16} m^2 e^2 \gamma^2 \quad \text{,,} \quad 2c - 2g,$$

whence, writing as above $l + \delta l$ for l , and integrating, we have

$$\delta l = -\frac{5}{16} e^2 \gamma^2 \sin \quad 2c - 2g,$$

and it thus appears that the values of δa , δe , $\delta \gamma$, δc , δg , δl , agree with those obtained in my former Note.