

44.

ON A MULTIPLE INTEGRAL CONNECTED WITH THE THEORY OF ATTRACTIONS.

[From the *Cambridge and Dublin Mathematical Journal*, vol. II. (1847), pp. 219—223.]

MR BOOLE [in the Memoir "On a Certain Multiple definite Integral" Irish Acad. Trans. vol. XXI. (1848), pp. 40—150] has given for the integral with n variables

$$V = \int \frac{\phi \left(\frac{x^2}{f^2} + \frac{y^2}{g^2} + \dots \right) dx dy \dots}{[(a-x)^2 + (b-y)^2 \dots + u^2]^{\frac{1}{2}n+q}} \dots \dots \dots (1);$$

limits
$$\frac{x^2}{f^2} + \frac{y^2}{g^2} + \dots \leq 1,$$

the following formula, or one which may readily be reduced to that form¹,

$$V = \frac{fg \dots \pi^{\frac{1}{2}n}}{\Gamma(\frac{1}{2}n + q)} \int_{\eta}^{\infty} \frac{Ss^{-q-1} ds}{\sqrt{\{(s+f^2)(s+g^2)\dots\}}} \dots \dots \dots (2),$$

where

$$S = \frac{(1-\sigma)^{-q}}{\Gamma(-q)} \int_0^1 t^{-q-1} \phi \{ \sigma + t(1-\sigma) \} dt \dots \dots \dots (3);$$

in which

$$\sigma = \frac{a^2}{f^2+s} + \frac{b^2}{g^2+s} \dots + \frac{u^2}{s} \dots \dots \dots (4)$$

and η is determined by

$$1 = \frac{a^2}{f^2+\eta} + \frac{b^2}{g^2+\eta} \dots + \frac{u^2}{\eta}.$$

¹ See note at the end of this paper.

Suppose $f = g = \dots = \infty$; also assume

$$\phi(\lambda) = \frac{1}{(f^2\lambda + v^2)^{\frac{1}{2}n+q}} \dots\dots\dots(5);$$

then the integral becomes

$$U = \int \frac{dx dy \dots}{(x^2 + y^2 \dots + v^2)^{\frac{1}{2}n+q} \{(x-a)^2 + (y-b)^2 \dots + u^2\}^{\frac{1}{2}n+q}} \dots\dots\dots(6).$$

the limits for each variable being $-\infty, \infty$.

Now, writing f^2s for s and $f^2\eta$ for η , the new value of η reduces itself to zero, and

$$U = \frac{f^{-2q} \pi^{\frac{1}{2}n}}{\Gamma(\frac{1}{2}n + q)} \int_0^\infty \frac{S ds}{(1+s)^{\frac{1}{2}n}}.$$

Also $\sigma = 0$; but

$$f^2\sigma = \frac{r^2}{1+s} + \frac{u^2}{s},$$

where $r^2 = a^2 + b^2 + \dots$ whence also putting $\frac{t}{f^2}$ for t , $\phi\{\sigma + t(1-\sigma)\}$ becomes

$$\frac{1}{\{f^2\sigma + t(1-\sigma) + v^2\}^{\frac{1}{2}n+q}};$$

i.e.

$$\frac{1}{(t+A)^{\frac{1}{2}n+q}},$$

if for a moment

$$A = \frac{r^2}{1+s} + \frac{u^2}{s} + v^2.$$

Hence

$$\begin{aligned} S &= \frac{f^{2q}}{\Gamma(-q)} \int_0^\infty \frac{t^{-q-1} dt}{(t+A)^{\frac{1}{2}n+q}} \\ &= \frac{\Gamma(\frac{1}{2}n + q + q')}{\Gamma(\frac{1}{2}n + q')} \frac{f^{2q}}{A^{\frac{1}{2}n+q+q'}}; \end{aligned}$$

or substituting in U , and replacing A by its value,

$$U = \frac{\pi^{\frac{1}{2}n} \Gamma(\frac{1}{2}n + q + q')}{\Gamma(\frac{1}{2}n + q) \Gamma(\frac{1}{2}n + q')} \int_0^\infty \frac{s^{-q-1} ds}{(1+s)^{\frac{1}{2}n} \left(\frac{r^2}{1+s} + \frac{u^2}{s} + v^2\right)^{\frac{1}{2}n+q+q'}};$$

or, what comes to the same,

$$U = \frac{\pi^{\frac{1}{2}n} \Gamma(\frac{1}{2}n + q + q')}{\Gamma(\frac{1}{2}n + q) \Gamma(\frac{1}{2}n + q')} \int_0^\infty \frac{s^{\frac{1}{2}n+q'-1} (1+s)^{q+q'} ds}{(v^2s^2 + js + u^2)^{\frac{1}{2}n+q+q'}} \dots\dots\dots(7)$$

where

$$j = u^2 + v^2 + r^2.$$

The only practicable case is that of $q' = -q$, for which

$$U = \frac{\pi^{\frac{1}{2}n} \Gamma(\frac{1}{2}n)}{\Gamma(\frac{1}{2}n + q) \Gamma(\frac{1}{2}n - q)} \int_0^\infty \frac{s^{\frac{1}{2}n - q - 1} ds}{(v^2 s^2 + js + u^2)^{\frac{1}{2}n}} \dots\dots\dots(8).$$

Consider the more general expression

$$\Theta = \int_0^\infty s^{-q-1} \phi\left(\frac{v^2 s^2 + js + u^2}{s}\right) ds \dots\dots\dots(9);$$

by writing

$$2u\sqrt{s} = \sqrt{(s' + 4uv)} \pm \sqrt{s'},$$

the upper sign from $s = \infty$ to $s = \frac{u}{v}$, and the lower one from $s = \frac{u}{v}$ to $s = 0$, it is easy to derive

$$\Theta = (2v)^{2q} \int_0^\infty \frac{\{\sqrt{(s + 4uv)} + \sqrt{s}\}^{-2q} + \{\sqrt{(s + 4uv)} - \sqrt{s}\}^{-2q}}{\sqrt{s} \sqrt{(s + 4uv)}} \phi(s + j + 2uv) ds \dots\dots\dots(10).$$

Now, by a formula which will presently be demonstrated,

$$\begin{aligned} & \int_0^\infty \frac{\{\sqrt{(s + 4uv)} + \sqrt{s}\}^{-2q} + \{\sqrt{(s + 4uv)} - \sqrt{s}\}^{-2q}}{\sqrt{s} \sqrt{(s + 4uv)}} e^{-\theta s} ds \\ &= \frac{2\sqrt{\pi}}{\Gamma(\frac{1}{2} - q)} \theta^{-q} \int_0^\infty s^{-\frac{1}{2} - q} (s + 4uv)^{-\frac{1}{2} - q} e^{-\theta s} ds \dots\dots\dots(11); \end{aligned}$$

whence

$$\begin{aligned} & \int_0^\infty \frac{\{\sqrt{(s + 4uv)} + \sqrt{s}\}^{-2q} + \{\sqrt{(s + 4uv)} - \sqrt{s}\}^{-2q}}{\sqrt{s} \sqrt{(s + 4uv)}} F_s . ds \\ &= \frac{2\sqrt{\pi}}{\Gamma(\frac{1}{2} - q)} \int_0^\infty s^{-\frac{1}{2} - q} (s + 4uv)^{-\frac{1}{2} - q} \left(-\frac{d}{ds}\right)^{-q} F_s . ds \dots\dots\dots(12). \end{aligned}$$

Thus, by merely changing the function,

$$\Theta = \frac{2^{2q+1} v^{2q} \sqrt{\pi}}{\Gamma(\frac{1}{2} - q)} \int_0^\infty s^{-\frac{1}{2} - q} (s + 4uv)^{-\frac{1}{2} - q} \left(-\frac{d}{ds}\right)^{-q} \phi(s + j + 2uv) ds \dots\dots\dots(13);$$

and hence in the particular case in question

$$U = \frac{2^{2q+1} v^{2q} \pi^{\frac{1}{2}(n+1)}}{\Gamma(\frac{1}{2} - q) \Gamma(\frac{1}{2}n + q)} \int_0^\infty s^{-\frac{1}{2} - q} (s + 4uv)^{-\frac{1}{2} - q} (s + j + 2uv)^{-\frac{1}{2}n + q} ds \dots\dots\dots(14),$$

by means of the formula

$$\left(-\frac{d}{ds}\right)^{-q} (s + \alpha)^{-\frac{1}{2}n} = \frac{\Gamma(\frac{1}{2}n - q)}{\Gamma(\frac{1}{2}n)} (s + \alpha)^{-\frac{1}{2}n + q}.$$

But as there may be some doubt about this formula, which is not exactly equivalent either to Liouville's or Peacock's expression for the general differential coefficient of a

power, it is worth while to remark that, by first transforming the $\frac{1}{2}n^{\text{th}}$ power into an exponential, and then reducing as above (thus avoiding the general differentiation), we should have obtained

$$U = \frac{2^{2q+1} v^{2q} \pi^{\frac{1}{2}(n+1)}}{\Gamma(\frac{1}{2}-q) \Gamma(\frac{1}{2}n+q) \Gamma(\frac{1}{2}n-q)} \int_0^\infty d\theta \int_0^\infty ds \theta^{\frac{1}{2}n-q-1} e^{-\theta(s+j+2uv)} s^{-\frac{1}{2}-q} (s+4uv)^{-\frac{1}{2}-q} e^{-\theta s},$$

which reduces itself to the equation (14) by simply performing the integration with respect to θ ; thus establishing the formula beyond doubt¹. The integral may evidently be effected in finite terms when either q or $q - \frac{1}{2}$ is integral. Thus for instance in the simplest case of all, or when $q = -\frac{1}{2}$,

$$U = \frac{\pi^{\frac{1}{2}(n-1)}}{v \Gamma(\frac{1}{2}(n+1))} \frac{1}{(j+2uv)^{\frac{1}{2}(n-1)}} = \int_{-\infty}^\infty \frac{dx dy \dots}{(x^2 + y^2 \dots + v^2)^{\frac{1}{2}(n+1)} \{(x-a)^2 + \dots u^2\}^{\frac{1}{2}(n-1)}},$$

a formula of which several demonstrations have already been given in the *Journal*.

The following is a demonstration, though an indirect one, of the formula (11): in the first place

$$\begin{aligned} & \int_0^\infty \frac{\{\sqrt{(s+4uv)} + \sqrt{s}\}^{-2q} + \{\sqrt{(s+4uv)} - \sqrt{s}\}^{-2q}}{\sqrt{s} \sqrt{(s+4uv)}} e^{-\theta s} ds \\ &= \frac{2\Gamma(\frac{1}{2}-q)}{\sqrt{\pi} (4uv)^{2q}} \theta^q e^{2uv\theta} \int_0^\infty (4u^2v^2 + x^2)^{q-\frac{1}{2}} e^{i\theta x} dx \dots\dots\dots(16), \end{aligned}$$

(where as usual $i \doteq \sqrt{-1}$): to prove this, we have

$$\begin{aligned} \int_{-\infty}^\infty (4u^2v^2 + x^2)^{q-\frac{1}{2}} e^{i\theta x} dx &= \frac{1}{\Gamma(\frac{1}{2}-q)} \int_{-\infty}^\infty dx \int_0^\infty dt t^{-\frac{1}{2}-q} e^{-t(4u^2v^2+x^2)+i\theta x} \\ &= \frac{\sqrt{\pi}}{\Gamma(\frac{1}{2}-q)} \int_0^\infty dt t^{-1-q} e^{-4u^2v^2t - \frac{\theta^2}{4t}}; \end{aligned}$$

or, putting $4uv \sqrt{t} = \sqrt{(s+4uv)} \pm \sqrt{s}$ (which is a transformation already employed in the present paper), the formula required follows immediately.

Now, by a formula due to M. Catalan, but first rigorously demonstrated by M. Serret,

$$\int_0^\infty \frac{\cos \alpha x dx}{(1+x^2)^n} = \frac{\pi}{(\Gamma n)^2} \int_0^\infty e^{-(\alpha+2z)} (z+\alpha)^{n-1} z^{n-1} dz,$$

(*Liouville*, t. VIII. [1843] p. 1), and by a slight modification in the form of this equation

$$\int_{-\infty}^\infty (4u^2v^2 + x^2)^{q-\frac{1}{2}} e^{i\theta x} dx = \frac{\pi e^{-2uv\theta} (4uv)^{2q}}{\theta^{2q} \Gamma^2(\frac{1}{2}-q)} \int_0^\infty s^{-q-\frac{1}{2}} (s+4uv)^{-q-\frac{1}{2}} e^{-\theta s} ds,$$

which, compared with (16), gives the required equation.

¹ A paper by M. Schlömilch "Note sur la variation des constantes arbitraires d'une Integrale definie," *Crelle*, t. xxxiii. [1846], pp. 268—280, will be found to contain formulæ analogous to some of the preceding ones.

NOTE.—One of the intermediate formulæ of Mr Boole [in the Memoir referred to] may be written as follows:

$$S = \frac{1}{\pi} \int_0^1 d\alpha \int_0^\infty dv v^q \cos [(\alpha - \sigma)v + \frac{1}{2}q\pi] \phi\alpha,$$

or what comes to the same thing, putting $i = \sqrt{-1}$, and rejecting the impossible part of the integral,

$$S = \frac{1}{\pi} e^{\frac{1}{2}q\pi i} \int_0^1 d\alpha \int_0^\infty dv v^q e^{2v(\alpha - \sigma)} \phi\alpha,$$

i.e.

$$S = \int_0^1 I \phi\alpha d\alpha, \quad I = \frac{1}{\pi} e^{\frac{1}{2}q\pi i} \int_0^\infty dv v^q e^{iv(\alpha - \sigma)}.$$

Now $(\alpha - \sigma)$ being positive,

$$I = \frac{1}{\pi} e^{\frac{1}{2}q\pi i} \Gamma(q+1) e^{\frac{1}{2}(q+1)\pi i} (\alpha - \sigma)^{-q-1};$$

i.e.

$$I = \frac{1}{\pi} e^{(q+\frac{1}{2})\pi i} \Gamma(q+1) (\alpha - \sigma)^{-q-1},$$

or, retaining the real part only,

$$I = -\frac{1}{\pi} \sin q\pi \Gamma(q+1) (\alpha - \sigma)^{-q-1};$$

i.e.

$$I = \frac{1}{\Gamma(-q)} (\alpha - \sigma)^{-q-1}.$$

But $(\alpha - \sigma)$ being negative,

$$I = \frac{1}{\pi} e^{\frac{1}{2}q\pi i} \Gamma(q+1) e^{-\frac{1}{2}(q+1)\pi i} (\sigma - \alpha)^{-q-1};$$

i.e.

$$I = \frac{1}{\pi} e^{-\frac{1}{2}\pi i} \Gamma(q+1) (\sigma - \alpha)^{-q-1},$$

or, retaining the real part only, $I = 0$.

Hence

$$S = \frac{1}{\Gamma(-q)} \int_\sigma^1 (\alpha - \sigma)^{-q-1} \phi\alpha d\alpha;$$

or putting

$$\alpha = \sigma + t(1 - \sigma), \text{ or } \alpha - \sigma = t(1 - \sigma),$$

$$S = \frac{(1 - \sigma)^{-q}}{\Gamma(-q)} \int_0^1 t^{-q-1} \phi[\sigma + t(1 - \sigma)] dt;$$

the expression in the text. Mr Boole's final value is

$$S = \left(-\frac{d}{d\sigma}\right)^q \phi(\sigma),$$

which, though simpler, appears to me to be in some respects less convenient.