

## 484.

ON THE VARIATIONS OF THE POSITION OF THE ORBIT IN  
THE PLANETARY THEORY.

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It has always appeared to me that in the Planetary Theory, more especially when the method of the variation of the elements is made use of, there is a difficulty as to the proper mode of dealing with the inclinations and longitudes of the nodes, hindering the ulterior development of the theory. Considering the case of two planets  $m$ ,  $m'$ , and referring their orbits to any fixed plane and fixed origin of longitudes therein, let  $\theta$ ,  $\theta'$  be the longitudes of the nodes,  $\phi$ ,  $\phi'$  the inclinations ( $p = \tan \phi \sin \theta$ ,  $q = \tan \phi \cos \theta$ , &c., as usual); then the disturbing functions for  $m$ ,  $m'$  respectively are developed, not explicitly in terms of  $\phi$ ,  $\phi'$ ,  $\theta$ ,  $\theta'$ , but in terms of  $\Phi$ , the mutual inclination of the two orbits, and of  $\Theta$ ,  $\Theta'$  the longitudes in the two orbits respectively of the mutual node of the two orbits;  $\Phi$  and  $\Theta$ ,  $\Theta'$  being functions (and complicated ones) of  $\phi$ ,  $\phi'$ ,  $\theta$ ,  $\theta'$ . Moreover, although in the general theory of the secular variations of the orbits of the planetary system,  $\theta$ ,  $\phi$ , &c., are, as above, referred to one fixed plane (the ecliptic of a certain date), yet in the theory of each particular planet it is the practice, and obviously the convenient one, to refer for such planet the  $\theta$ ,  $\phi$  to its own fixed plane (the orbit of the planet at a certain date), the effect of course being that  $\phi$ , and consequently  $p$ ,  $q$ , instead of being of the order of the inclinations to the ecliptic, are only of the order of the disturbing forces. It has occurred to me that the last-mentioned plan should be adhered to *throughout*; viz., that for each planet  $m$ , the position of its variable orbit should be determined by  $\theta$ , the longitude of its node, and  $\phi$ , the inclination in reference to the appropriate fixed plane (orbit of the planet at a certain date) and origin of longitude therein. The disturbing functions for the planets  $m$  and  $m'$  will of course depend not only on  $\theta$ ,  $\theta'$ ,  $\phi$ ,  $\phi'$ , but on the quantities  $\Phi$ ,  $\Theta$ ,  $\Theta'$  which determine the mutual positions of the two fixed

planes of reference and origins of longitude therein, *these last being however absolute constants not affected by any variation of the elements*; so that as regards the variation of the elements the disturbing functions are in fact given as *explicit* functions of the variable elements  $\theta, \theta', \phi, \phi'$ ; and where  $\phi, \phi'$  and therefore also  $p, q, p', q'$  are only of the order of the disturbing forces.

I proceed to work out this idea, for the present considering the development of the Disturbing Function only as far as the first powers of  $p, q, \&c.$  For comparison with the ordinary theory, observe that in this theory the disturbing function contains only the *second* powers of the  $p, q, \&c.$ , made use of therein; these are in fact of a form such as  $P+p, Q+q, \dots$  where  $P, Q$  are absolute constants and  $p, q, \dots$  are the  $p, q, \dots$  of the present theory; the ordinary theory gives therefore in the disturbing function a series of terms involving  $(P+p)^2, (P+p)(Q+q), \dots$  which I now take account of only as far as the first powers of  $p, q, \dots$  viz., they are in effect reduced to  $P^2+2Pp, PQ+Pq+Qp, \&c. \dots$  The present theory is thus not now developed to the extent of giving the  $p, q, \dots$  of the ordinary theory in the more complete form as the solutions of a system of simultaneous linear differential equations, but only to the extent of obtaining for these  $p, q, \dots$  respectively the terms which are proportional to the time.

I commence with the following subsidiary problem. Consider a spherical triangle  $ABC$  (sides  $a, b, c$ , angles  $A, B, C$ , as usual), and taking the side  $c$  as constant, but the angles  $A$  and  $B$  as variable, let it be required to find the variations of  $C, a, b$  in terms of variations  $dA, dB$  and the variable elements  $C, a, b$  themselves. Although the geometrical proof would be more simple, I give the analytical one, as it may be useful.

We have

$$\cos C = -\cos A \cos B + \sin A \sin B \cos c,$$

and thence

$$\begin{aligned} -\sin C dC &= (\sin A \cos B + \cos A \sin B \cos c) dA \\ &\quad + (\sin B \cos A + \sin A \cos B \cos c) dB \\ &= \frac{\sin B \sin c}{\tan b} dA + \frac{\sin A \sin c}{\tan a} dB, \end{aligned}$$

that is

$$-\frac{\sin C}{\sin c} dC = \frac{\sin B \cos b}{\sin b} dA + \frac{\sin A \cos a}{\sin a} dB,$$

or finally

$$-dC = \cos b dA + \cos a dB.$$

Next

$$\sin a = \sin c \frac{\sin A}{\sin C},$$

or, differentiating,

$$\cos a da = \frac{\sin c}{\sin^2 C} (\sin C \cos A dA - \cos C \sin A dC)$$

or, substituting for  $dC$  its value,

$$\begin{aligned}
 &= \frac{\sin c}{\sin^2 C} \left\{ dA (\sin C \cos A + \cos C \sin A \cos b) + dB \cos C \sin A \cos a \right\}, \\
 &= \frac{\sin c}{\sin^2 C} \left\{ dA \frac{\sin A \sin b \cos a}{\sin a} + dB \cos C \sin A \cos a \right\},
 \end{aligned}$$

that is

$$da = \frac{1}{\sin C} \left\{ dA \frac{\sin A}{\sin a} \sin b + dB \cos C \sin A \right\} \div \frac{\sin C}{\sin c},$$

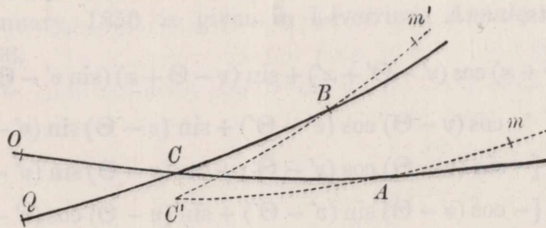
or, on the right-hand writing  $\frac{\sin A}{\sin a}$  instead of  $\frac{\sin C}{\sin c}$ , this is

$$da = \frac{1}{\sin C} (dA \sin b + dB \cos C \sin a);$$

and similarly

$$db = \frac{1}{\sin C} (dB \sin a + dA \cos C \sin b).$$

Now let the continuous lines represent the orbits of  $m, m'$  at certain dates,  $O, Q$  the origins of longitude therein; and the dotted lines the variable orbits of the planets respectively.



Write

$$\begin{aligned}
 OC &= \Theta, & CA &= \theta, & \angle CAC' &= \phi, \\
 QC &= \Theta', & CB &= \theta', & \angle CCB &= \phi', \\
 \angle C &= \Phi.
 \end{aligned}$$

Then, answering to the notation of the lemma, we have

$$\begin{aligned}
 a &= \theta', & b &= \theta, & C &= \Phi, & dA &= \phi, & dB &= -\phi', \\
 & & & & & & \text{or say} & = \tan \phi, & & = -\tan \phi',
 \end{aligned}$$

whence

$$\begin{aligned}
 C'B &= a + da, \\
 &= \theta' + \frac{1}{\sin \Phi} (\tan \phi \sin \theta - \tan \phi' \cos \Phi \sin \theta'), \\
 &= \theta' + \frac{1}{\sin \Phi} (p - p' \cos \Phi),
 \end{aligned}$$

$$\begin{aligned}
 C'A &= b + db, \\
 &= \theta + \frac{1}{\sin \Phi} (-\tan \theta' \sin \phi' + \tan \phi \cos \Phi \sin \theta), \\
 &= \theta - \frac{1}{\sin \Phi} (p' - p \cos \Phi), \\
 \angle C' &= C + dC = \Phi - \cos \theta \tan \phi + \cos \theta' \tan \phi', = \Phi - q + q'.
 \end{aligned}$$

Suppose  $v, v'$  are the longitudes of the planets in their two orbits respectively; that is

$$\begin{aligned}
 v &= OA + Am = \Theta + \theta + Am, \\
 v' &= QB + Bm' = \Theta + \theta' + Bm',
 \end{aligned}$$

whence

$$\begin{aligned}
 C'm &= C'A + Am, = v - \Theta - \frac{1}{\sin \Phi} (p' - p \cos \Phi), \\
 C'm' &= C'B + Bm, = v' - \Theta' + \frac{1}{\sin \Phi} (p - p' \cos \Phi), \\
 \angle C' &= \Phi - q + q';
 \end{aligned}$$

say these values are  $v - \Theta + x, v' - \Theta' + x', \Phi + y$ . Then if  $H$  is the angular distance  $mm'$  of the two planets,

$$\begin{aligned}
 \cos \bar{H} &= \cos (v - \Theta + x) \cos (v' - \Theta' + x') + \sin (v - \Theta + x) (\sin v' - \Theta' + x') \cos (\Phi + y), \\
 &= \cos (v - \Theta) \cos (v' - \Theta') + \sin (v - \Theta) \sin (v' - \Theta') \cos \Phi \\
 &\quad + x [-\sin (v - \Theta) \cos (v' - \Theta') + \cos (v - \Theta) \sin (v' - \Theta') \cos \Phi] \\
 &\quad + x' [-\cos (v - \Theta) \sin (v' - \Theta') + \sin (v - \Theta) \cos (v' - \Theta') \cos \Phi] \\
 &\quad + y [-\sin (v - \Theta) \sin (v' - \Theta') \sin \Phi], \\
 &= \cos H + \nabla \text{ suppose.}
 \end{aligned}$$

The disturbing function for the planet  $m$  disturbed by  $m'$  is

$$\Omega = m' \left\{ \frac{1}{\sqrt{r^2 + r'^2 - 2rr' \cos \bar{H}}} - \frac{r \cos \bar{H}}{r'^2} \right\},$$

( $\Omega = -R$ , if  $R$  is the disturbing function of the *Mécanique Céleste*); and the term hereof which involves  $\nabla$  is

$$= \nabla \frac{d\Omega}{d \cdot \cos H}$$

where after the differentiation  $\cos \bar{H}$  is replaced by  $\cos H$ ,

$$= m' \left\{ \frac{rr'}{(r^2 + r'^2 - 2rr' \cos H)^{\frac{3}{2}}} + \frac{r}{r'^2} \right\} \nabla,$$

viz., this is a linear function of  $x, x', y$ , that is of  $p, q, p', q'$ , with coefficients which of course involve the other variable elements and the time; but it will be remembered that  $\Theta, \Theta', \Phi$  are not variable elements, but are absolute constants. The variations of  $p$  depend upon  $\frac{d\Omega}{dq}$  and those of  $q$  on  $\frac{d\Omega}{dp}$ , and the quantities  $p, q, p', q', \dots$  disappear from these differential coefficients  $\frac{d\Omega}{dq}, \frac{d\Omega}{dp}$ ; that is, disregarding periodic terms, and the variations of the elements, we obtain  $\frac{dp}{dt}, \frac{dq}{dt}$  as absolute constants, or reckoning the time from the epoch belonging to the fixed orbit of  $m$ , we have  $p, q$  as mere multiples of the time ( $p = At, q = Bt$ , where  $A$  and  $B$  are constants); agreeing with the statement preceding the investigation.

Observe that the  $p, q$ , as used above, have reference not only to the fixed orbit of  $m$ , but also to the node thereon of the fixed orbit of  $m'$ : we may, if we please, write  $p = \tan \phi \sin (\Theta + \theta), q = \tan \phi \cos (\Theta + \theta)$ , that is,  $p = q \sin \Theta + p \cos \Theta, Q = q \cos \Theta - p \sin \Theta$  (or  $p = p \cos \Theta - q \sin \Theta, q = P \sin \Theta + q \cos \Theta$ ), and in place of  $p, q$  introduce into the formulæ  $p$  and  $q$ , which have reference only to the fixed orbit of  $m$ , and similarly writing  $p' = \tan \phi' \sin (\Theta' + \theta'), q' = \tan \phi' \cos (\Theta' + \theta')$ , instead of  $p', q'$  introduce  $p', q'$  which have reference only to the fixed orbit of  $m'$ .

I remark that a table for the relative positions of the orbits of the eight Planets for the Epoch 1st January, 1850, is given in Leverrier's *Annales de l'Observ. de Paris*, t. II. (1856), pp. 64—66.